

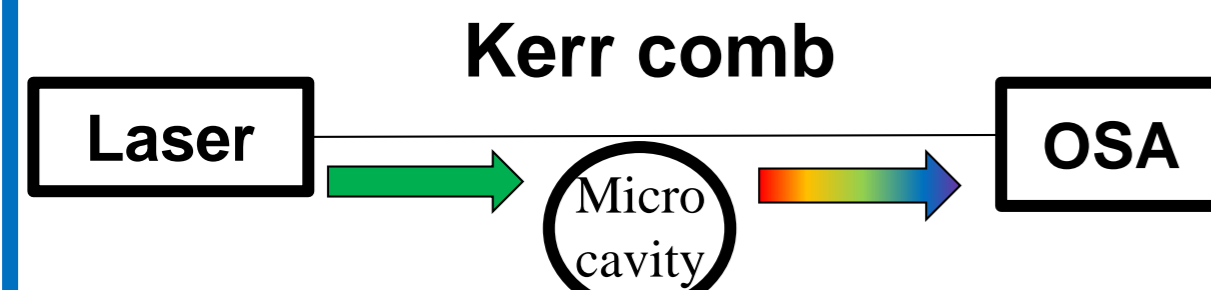
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Abstract

We numerically studied the effect of the interaction between transverse modes on Kerr frequency comb generation. We expanded the Lugiato-Lefever mode to consider the effect of cross-phase modulation (XPM). We found that XPM can work for forming solitons in both modes, TE and TM modes. This means dual solitons can be achieved in one microcavity.

Background: Kerr frequency comb generation

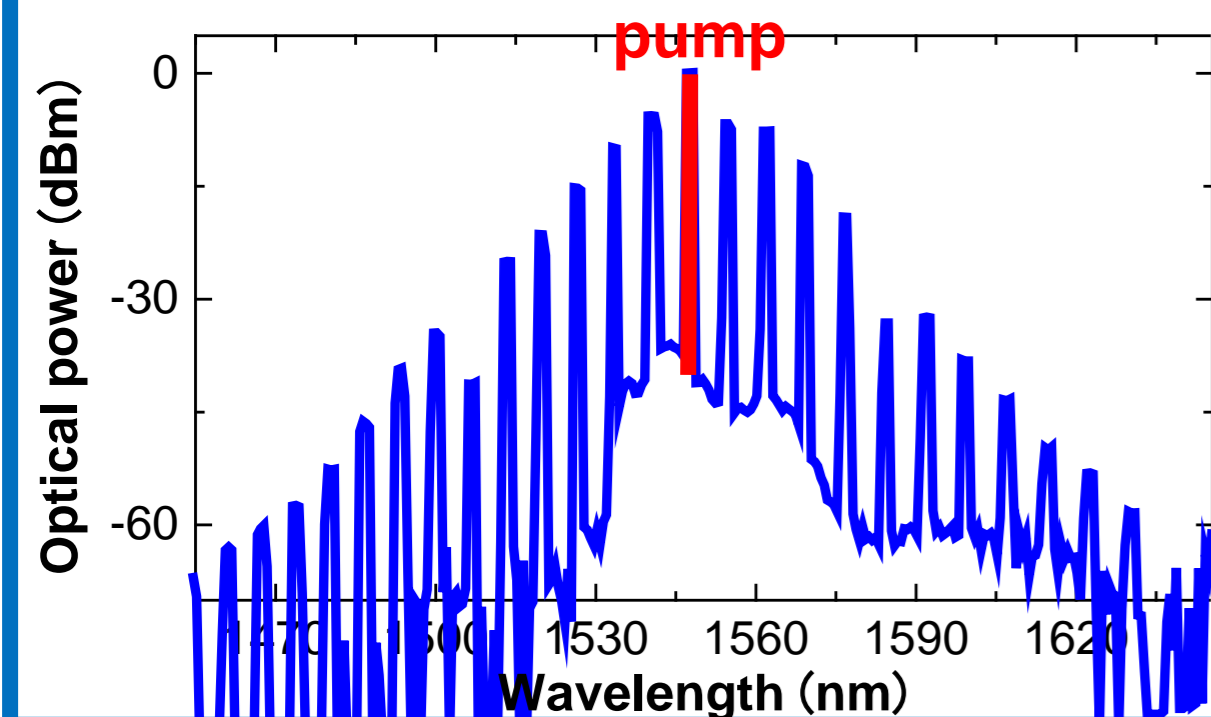


Kerr comb

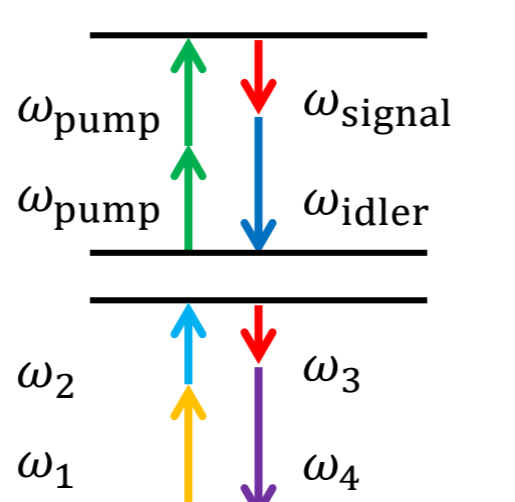
High repetition rate (> 100 GHz)
Small size (< 1 mm)
Low power consumption

Temporal soliton generation

- Mode-locking mechanism is clear.
- Stabilization has been developed.
- Some material platforms (MgF₂, SiN, Silica...)
- Variable generation band (telecom, Mid-IR, Vis)



Optical power (dBm) vs Wavelength (nm)



ω_{pump} ω_{signal}
 ω_{pump} ω_{idler}
 ω_2 ω_3
 ω_1 ω_4

Dual comb spectroscopy

- Two comb source
- Coherent link is needed.

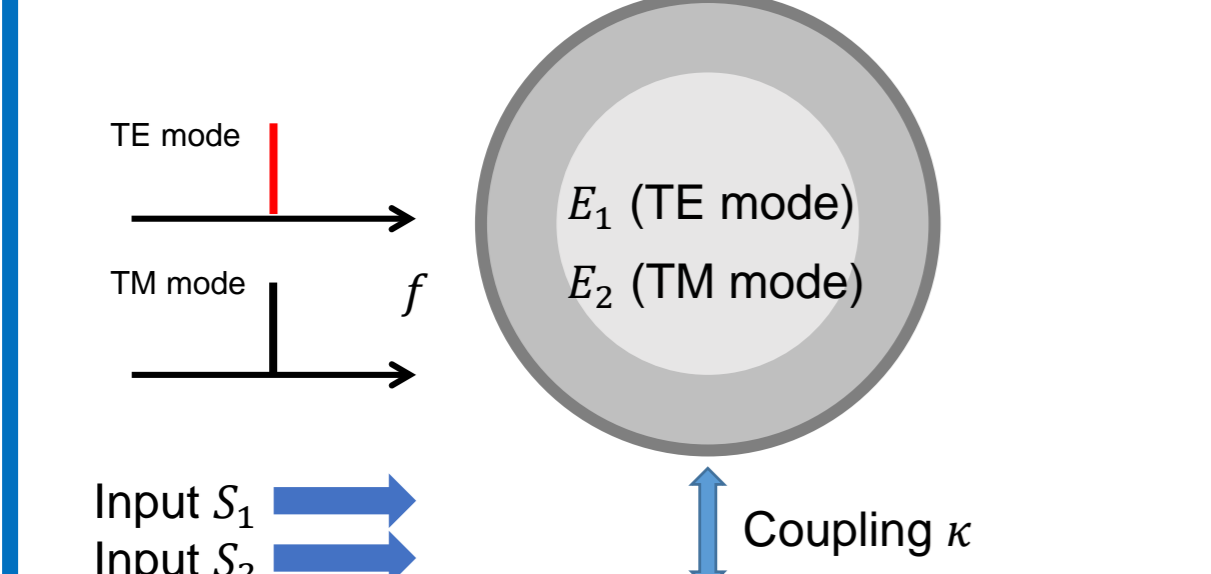
I. Coddington, *et al.*, *Optica* 3, 414 (2016).

For dual comb light source...

- An μ -cavity supports some high-Q mode families.
- An μ -cavity requires comparatively low input power.

Dual comb from dual input pumps could be achieved in one μ -cavity considering the interactions such as XPM.

Simulation Method: Transverse modes interaction



TE mode
TM mode
Input S_1
Input S_2
Coupling κ

S : Input power
 r : Round trip number
 t_R : Round trip time
 δ : Detuning
 M : Mode number

$$t_R \frac{\partial E_{TE}}{\partial r} = \left(-\alpha_{TE} - i\delta_{TE} + iL \sum_{k \geq 2} \frac{\beta_{k,TE}}{k!} \left(i \frac{\partial}{\partial T} \right)^k \right) E_{TE} + i\gamma_{TE} L (|E_{TE}|^2 + P|E_{TM}|^2) E_{TE} + \sqrt{\kappa_{TE}} S_{TE}$$

$$t_R \frac{\partial E_{TM}}{\partial r} = \left(-\alpha_{TM} - i\delta_{TM} + dLi \frac{\partial}{\partial T} + iL \sum_{k \geq 2} \frac{\beta_{k,TM}}{k!} \left(i \frac{\partial}{\partial T} \right)^k \right) E_{TM} + i\gamma_{TM} L (|E_{TM}|^2 + P|E_{TE}|^2) E_{TM} + \sqrt{\kappa_{TM}} S_{TM}$$

• Calculate them alternatively with split-step Fourier method

Mode interaction(XPM) coefficient P

- same polarization $P=2$
- orthogonal polarization $P=2/3$

For example,

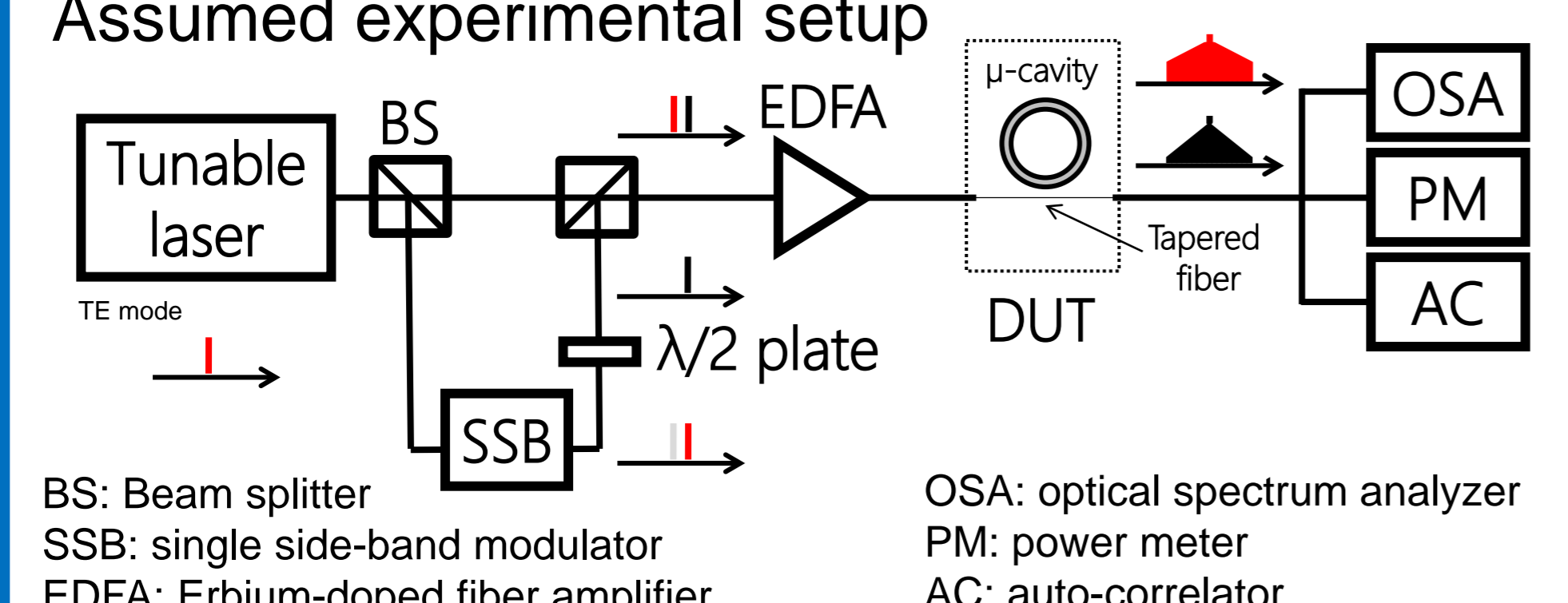
- 1st order TE \times 1st order TM: $P=2/3$
- 1st order TE \times 2nd order TE: $P=2$

To be simple,

- Mode overlapping is perfect. ($B=1$)
- Group velocity mismatch is negligible. ($d=0$)
- Linear coupling is negligible (only XPM is the interaction)

Simulation Result

Assumed experimental setup



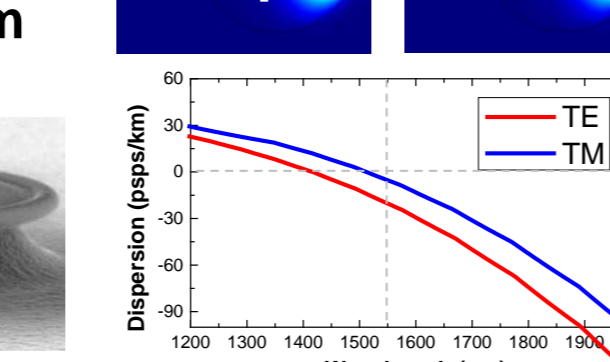
BS: Beam splitter
SSB: single side-band modulator
EDFA: Erbium-doped fiber amplifier

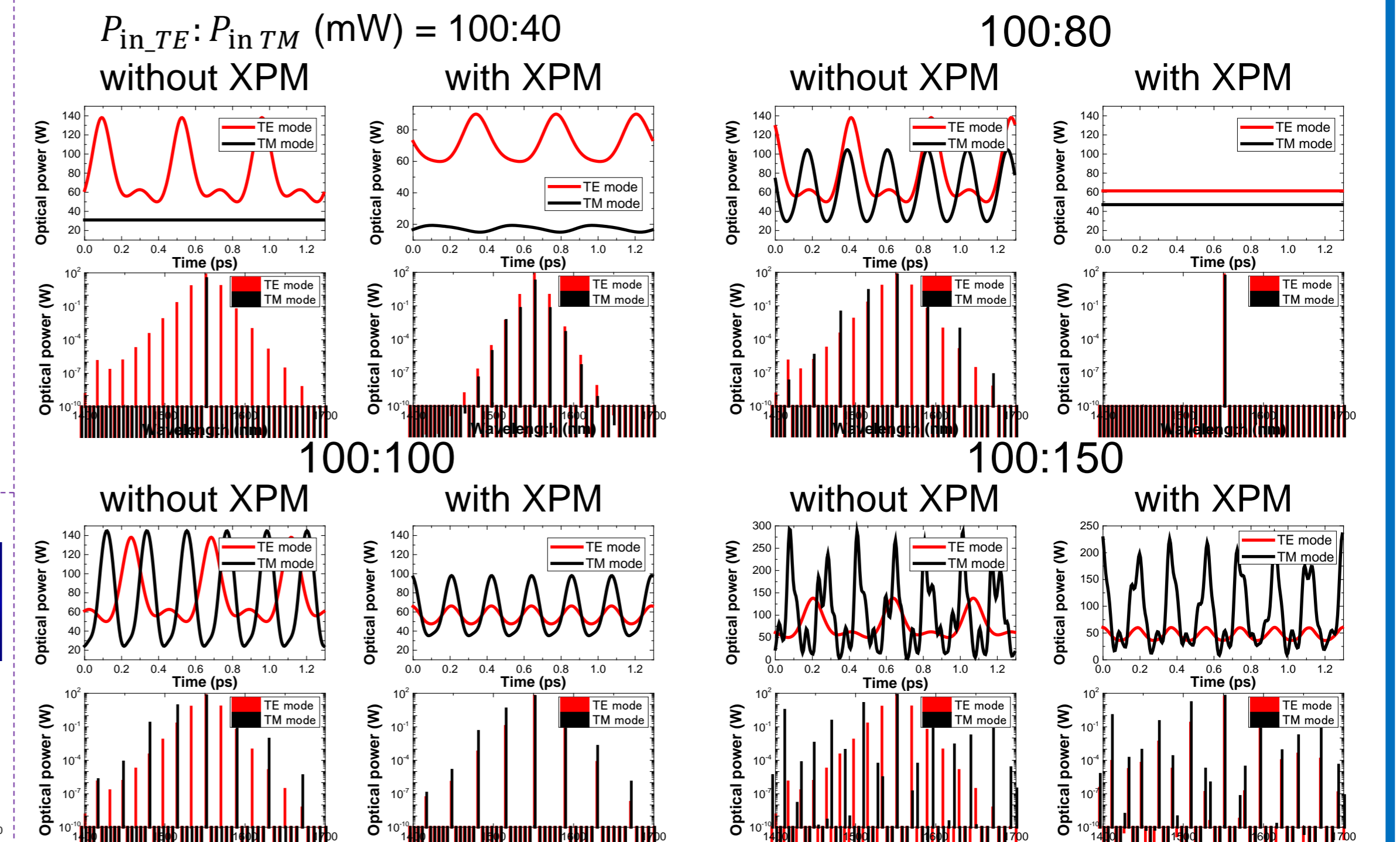
OSA: optical spectrum analyzer
PM: power meter
AC: auto-correlator

Calculation parameters

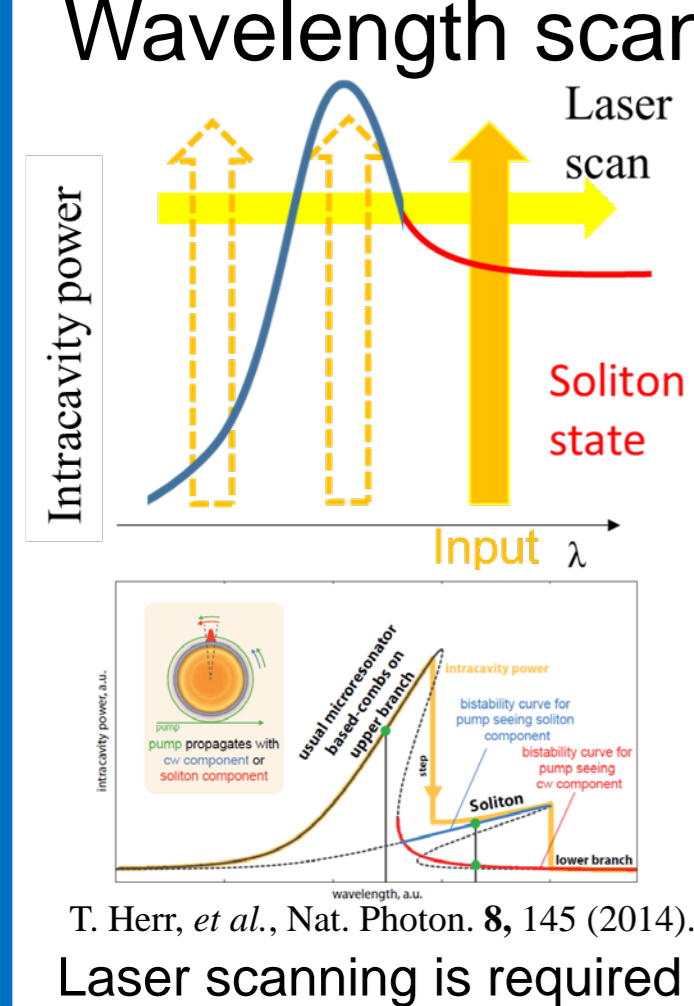
Q_{TE}	4.0×10^6
Q_{TM}	4.0×10^6
$A_{\text{eff-TE}}$	$4.63 \mu\text{m}^2$
$A_{\text{eff-TM}}$	$4.96 \mu\text{m}^2$
n_2	$2.2 \times 10^{-10} \text{ 1/Wm}$

Silica toroid
Diameter: 60 μm
FSR: 770 GHz





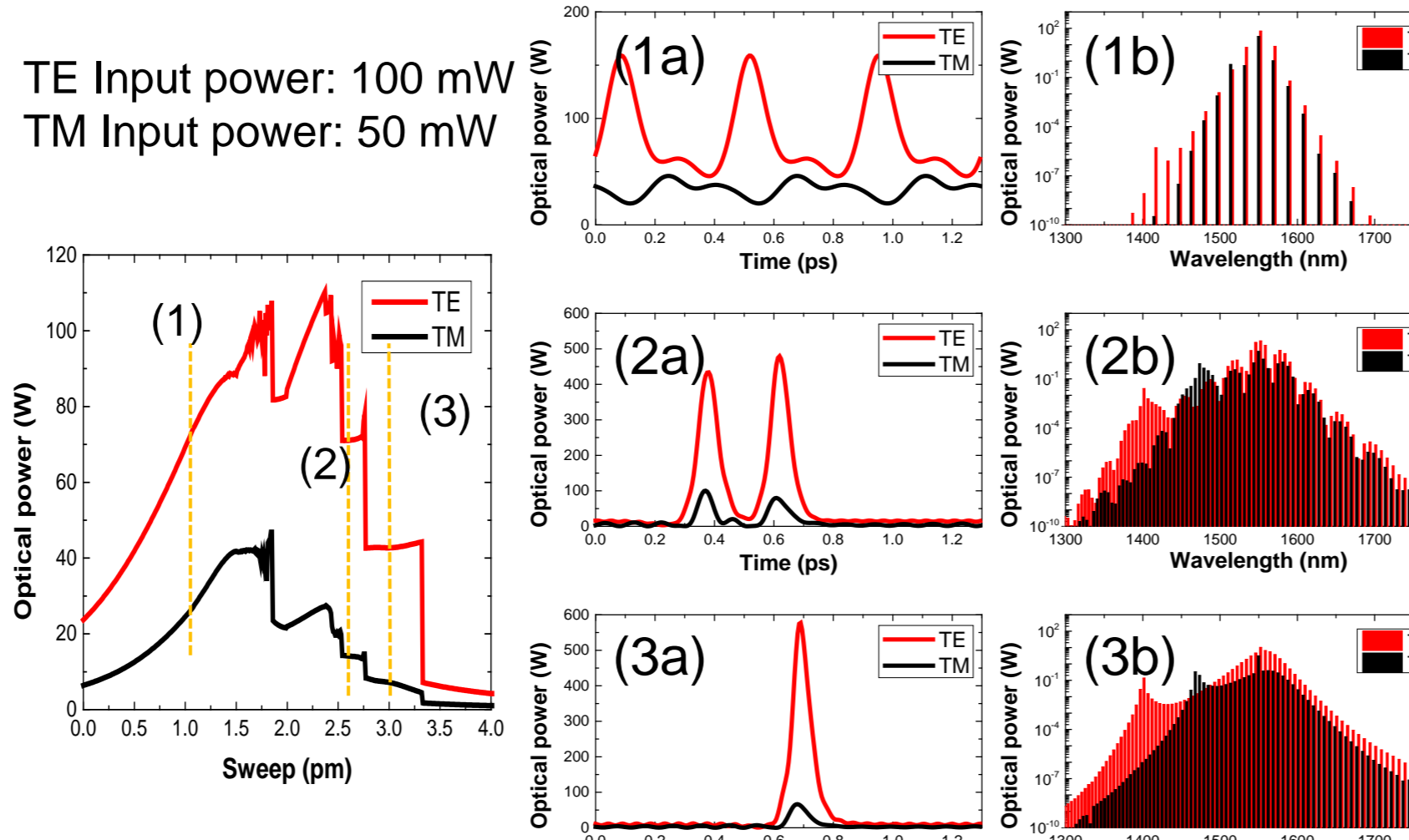
Wavelength scan



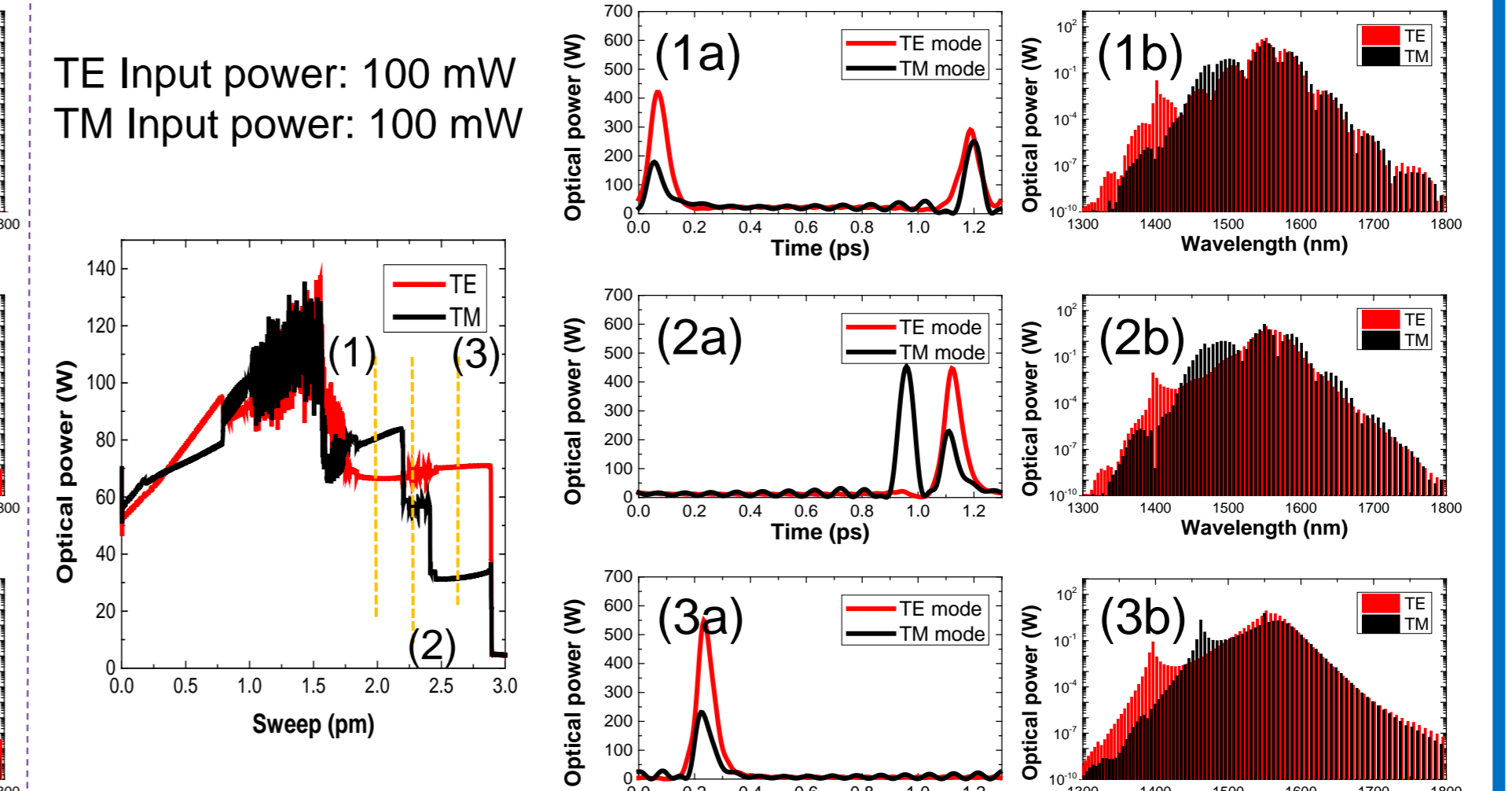
Laser scan
Soliton state
Input λ

T. Herr, *et al.*, *Nat. Photon.* 8, 145 (2014).
Laser scanning is required to reach the soliton state.

TE Input power: 100 mW
TM Input power: 50 mW



TE Input power: 100 mW
TM Input power: 100 mW



Conclusion

- ✓ Modelled the XPM effect with the SSFM calculation.
- ✓ Twin mode-locked pulses can be achieved with wavelength scanning method.
- ✓ Twin mode-locked pulses move forward at the same speed due to XPM effect that works like soliton trapping.

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